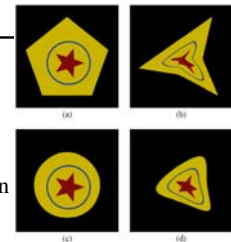
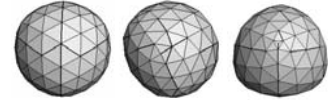


Mass-Spring Systems

Last Time?



- Subdivision Surfaces
 - Catmull Clark
 - Semi-sharp creases
 - Texture Interpolation
- Interpolation vs. Approximation

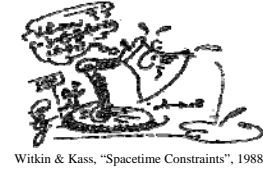


Today

- Particle Systems
 - Equations of Motion (Physics)
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Types of Dynamics

- Point 
- Rigid body 
- Deformable body (include clothes, fluids, smoke, etc.)



Witkin & Kass, "Spacetime Constraints", 1988.



Carlson, Mucha, Van Horn, & Turk 2002

What is a Particle System?

- Collection of many small simple particles that maintain *state* (position, velocity, color, etc.)
- Particle motion influenced by external *force fields*
- *Integrate* the laws of mechanics (ODE Solvers)
- To model: sand, dust, smoke, sparks, flame, water, etc.



Star Trek, The Wrath of Kahn 1982



Mark B. Allan
<http://users.rcn.com/mba.dnai/software/flow/>

Particle Motion

- mass m , position x , velocity v
- equations of motion:

$$\frac{d}{dt} x(t) = v(t)$$

$$\frac{d}{dt} v(t) = \frac{1}{m} F(x, v, t)$$

- Analytic solutions can be found for some classes of differential equations, but most can't be solved analytically
- Instead, we will numerically approximate a solution to our *initial value problem*

Path Through a Field

- $f(\mathbf{X}, t)$ is a vector field defined everywhere
 - E.g. a velocity field which may change over time



- $X(t)$ is a path through the field

Higher Order ODEs

- Basic mechanics is a 2nd order ODE:

$$\frac{d^2}{dt^2} x = \frac{1}{m} F$$

- Express as 1st order ODE by defining $v(t)$:

$$\frac{d}{dt} x(t) = v(t)$$

$$\frac{d}{dt} v(t) = \frac{1}{m} F(x, v, t)$$

$$\mathbf{X} = \begin{pmatrix} x \\ v \end{pmatrix} \quad f(\mathbf{X}, t) = \begin{pmatrix} v \\ \frac{1}{m} F(x, v, t) \end{pmatrix}$$

For a Collection of 3D particles...

$$\mathbf{X} = \begin{pmatrix} p_x^{(1)} \\ p_y^{(1)} \\ p_z^{(1)} \\ v_x^{(1)} \\ v_y^{(1)} \\ v_z^{(1)} \\ p_x^{(2)} \\ p_y^{(2)} \\ p_z^{(2)} \\ v_x^{(2)} \\ v_y^{(2)} \\ v_z^{(2)} \\ \vdots \end{pmatrix} \quad f(\mathbf{X}, t) = \begin{pmatrix} v_x^{(1)} \\ v_y^{(1)} \\ v_z^{(1)} \\ \frac{1}{m_1} F_x^{(1)}(\mathbf{X}, t) \\ \frac{1}{m_1} F_y^{(1)}(\mathbf{X}, t) \\ \frac{1}{m_1} F_z^{(1)}(\mathbf{X}, t) \\ v_x^{(2)} \\ v_y^{(2)} \\ v_z^{(2)} \\ \frac{1}{m_2} F_x^{(2)}(\mathbf{X}, t) \\ \frac{1}{m_2} F_y^{(2)}(\mathbf{X}, t) \\ \frac{1}{m_2} F_z^{(2)}(\mathbf{X}, t) \\ \vdots \end{pmatrix}$$

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Euler's Method

- Examine $f(\mathbf{X}, t)$ at (or near) current state
- Take a step of size h to new value of \mathbf{X} :

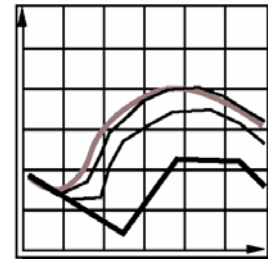
$$t_1 = t_0 + h$$

$$\mathbf{X}_1 = \mathbf{X}_0 + h f(\mathbf{X}_0, t_0)$$

- Piecewise-linear approximation to the curve

Effect of Step Size

- Step size controls accuracy
- Smaller steps more closely follow curve
- For animation, we may want to take many small steps per frame



Euler's Method: Inaccurate

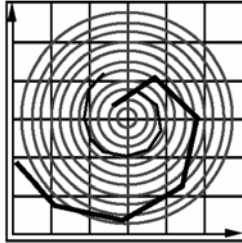
- Moves along tangent & can leave curve, e.g.:

$$f(\mathbf{X}, t) = \begin{pmatrix} -y \\ x \end{pmatrix}$$

- Exact solution is circle:

$$\mathbf{X}(t) = \begin{pmatrix} r \cos(t+k) \\ r \sin(t+k) \end{pmatrix}$$

- Euler's spirals outward no matter how small h is



Euler's Method: Unstable

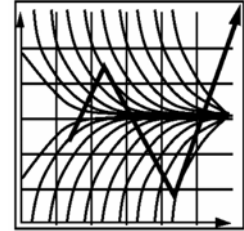
- Problem: $f(x, t) = -kx$
- Solution: $x(t) = x_0 e^{-kt}$

- Limited step size:

$$x_1 = x_0(1 - hk)$$

$$\begin{cases} h \leq 1/k & \text{ok} \\ h > 1/k & \text{oscillates } \pm \\ h > 2/k & \text{explodes} \end{cases}$$

- If k is big, h must be small



Analysis using Taylor Series

- Expand exact solution $\mathbf{X}(t)$

$$\mathbf{X}(t_0 + h) = \mathbf{X}(t_0) + h \left(\frac{d}{dt} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^2}{2!} \left(\frac{d^2}{dt^2} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^3}{3!} (\dots) + \dots$$

- Euler's method:

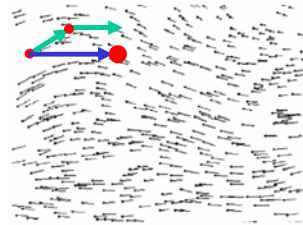
$$\mathbf{X}(t_0 + h) = \mathbf{X}_0 + h f(\mathbf{X}_0, t_0) \dots + O(h^2) \text{ error}$$

$$h \rightarrow h/2 \Rightarrow \text{error} \rightarrow \text{error}/4 \text{ per step} \times \text{twice as many steps} \rightarrow \text{error}/2$$

- First-order method: Accuracy varies with h
 - To get 100x better accuracy need 100x more steps

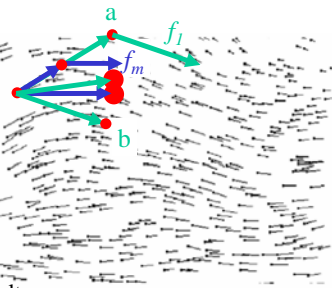
Can we do better than Euler's Method?

- Problem: f has varied along the step
- Idea: look at f at the arrival of the step and compensate for variation



2nd-Order Methods

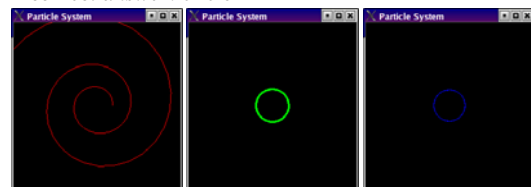
- Midpoint:**
 - 1/2 Euler step
 - evaluate f_m
 - full step using f_m
- Trapezoid:**
 - Euler step (a)
 - evaluate f_j
 - full step using f_j (b)
 - average (a) and (b)
- Not exactly same result
- Same order of accuracy



Comparison: Euler, Midpoint, Runge-Kutta

- initial position: (1,0,0)
- initial velocity: (0,5,0)
- force field: pulls particles to origin with magnitude proportional to distance from origin
- correct answer: circle

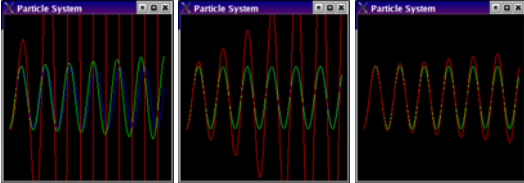
A 4th order method!



Euler will always diverge (even with small dt)

Comparison: Euler, Midpoint, Runge-Kutta

- *initial position*: (0,-2,0) A 4th order method!
- *initial velocity*: (1,0,0)
- *force field*: pulls particles to line y=0 with magnitude proportional to distance from line
- *correct answer*: sine wave



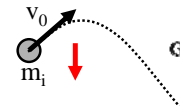
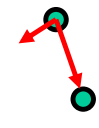
Decreasing the timestep (dt) improves the accuracy

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Forces: Gravity

For smoke, flame: make gravity point up!

- Simple gravity: depends only on particle mass
 - \mathbf{v}_0 
 - Gravity: $\mathbf{f}^{(i)} = \begin{pmatrix} 0 \\ 0 \\ -m_i G \end{pmatrix}$
- N-body problem: depends on all other particles
 - Magnitude inversely proportional to square distance
 - $F_{ij} = G m_i m_j / r^2$ 

Forces: Spatial Fields

- Force on particle i depends only on position of i
 - wind
 - attractors
 - repulsers
 - vortices
- Can depend on time
- Note: these add energy, may need damping too

Forces: Damping

$$\mathbf{f}^{(i)} = -d\mathbf{v}^{(i)}$$

- Force on particle i depends only on velocity of i
- Force opposes motion
- Removes energy, so system can settle
- Small amount of damping can stabilize solver
- Too much damping makes motion too glue-like

Questions?

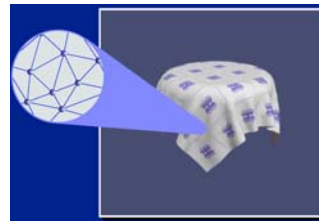


Image by Baraff, Witkin, Kass

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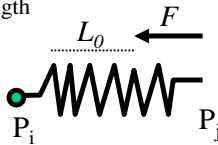
How would you simulate a string?

- Each particle is linked to two particles
- Forces try to keep the distance between particles constant
- What force?



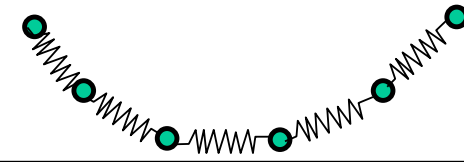
Spring Forces

- Force in the direction of the spring and proportional to difference with rest length L_0
- $$F(P_i, P_j) = K(L_0 - \|\vec{P}_i - \vec{P}_j\|) \frac{\vec{P}_i - \vec{P}_j}{\|\vec{P}_i - \vec{P}_j\|}$$
- K is the stiffness of the spring
 - When K gets bigger, the spring really wants to keep its rest length



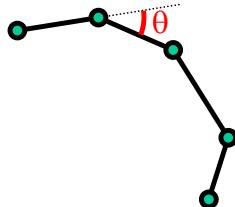
How would you simulate a string?

- Springs link the particles
- Springs try to keep their rest lengths and preserve the length of the string
- Problems?
 - Stretch, actual length will be greater than rest length
 - Numerical oscillation



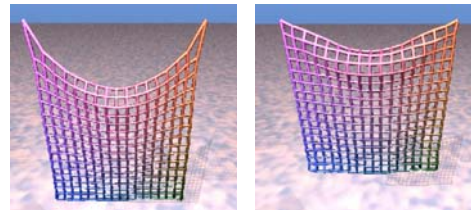
How would you simulate hair?

- Similar to string...
- Deformation forces proportional to the angle between segments



Reading for Today

- “Deformation Constraints in a Mass-Spring Model to Describe Rigid Cloth Behavior”, Provat, 1995.

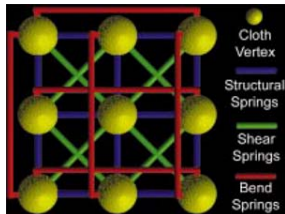


Simple mass-spring system

Improved solution

Cloth Modeled with Mass-Spring

- Network of masses and springs
- Structural springs:
 - link (i, j) & $(i+1, j)$
 - and (i, j) & $(i, j+1)$
- Shear springs
 - link (i, j) & $(i+1, j+1)$
 - and $(i+1, j)$ & $(i, j+1)$
- Flexion (Bend) springs
 - link (i, j) & $(i+2, j)$
 - and (i, j) & $(i, j+2)$



From Lander

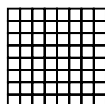
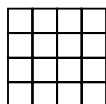
<http://www.darwin3d.com/gamesdev/articles/col0599.pdf>

The Stiffness Issue

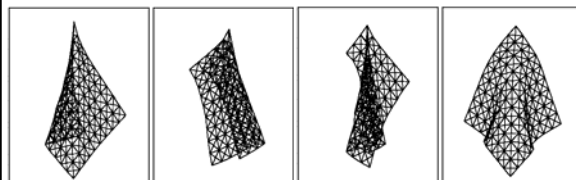
- What relative stiffness do we want for the different springs in the network?
- Cloth is barely elastic, shouldn't stretch so much!
- Inverse relationship between stiffness & Δt
- We really want a constraints (not springs)
- Many numerical solutions
 - reduce Δt
 - use constraints
 - implicit integration
 - ...

The Discretization Problem

- What happens if we discretize our cloth more finely, or with a different mesh structure?
- Do we get the same behavior?
- Usually not! It takes a lot of effort to design a scheme that does not depend on the discretization.



Questions?



Interactive Animation of Structured Deformable Objects
Desbrun, Schröder, & Barr 1999

Reading for Tuesday (2/5)

- Baraff, Witkin & Kass
Untangling Cloth, SIGGRAPH 2003

- Post a comment or question on the LMS discussion by 10am on Tuesday 2/5

