

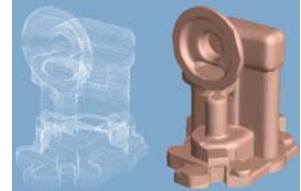
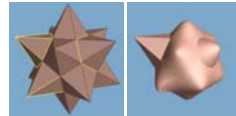
# Mass-Spring Systems

## Last Week?

- Spline Curves & Surfaces
- Subdivision Surfaces
  - Catmull Clark, Loop
  - Creases
  - Texture
  - Interpolation vs. Approximation
- Surface Reconstruction from Points





Figure 4: Interpolating a coarsely polygonized torus. Upper left: original mesh. Upper right: Spline surfaces interpolation. Lower left: Interpolating Catmull-Clark surface. Lower right: Patch interpolating Catmull-Clark surface.



## Today

- Particle Systems
  - Equations of Motion (Physics)
  - Numerical Integration (Euler, Midpoint, etc.)
  - Forces: Gravity, Spatial, Damping
- Mass Spring System Examples
  - String, Hair, Cloth
- Stiffness
- Discretization

## Types of Dynamics

- Point 
- Rigid body 
- Deformable body (include clothes, fluids, smoke, etc.)



Witkin & Kass, "Spacetime Constraints", 1988.



Carlson, Mucha, Van Horn, & Turk 2002

## What is a Particle System?

- Collection of many small simple particles that maintain *state* (position, velocity, color, etc.)
- Particle motion influenced by external *force fields*,
- *Integrate* the laws of mechanics (ODE Solvers)
- To model: sand, dust, smoke, sparks, flame, water, etc.

Star Trek, The Wrath of Kahn 1982



Mark B. Allan

<http://users.rcn.com/mba.dnai/software/flow/>

## Particle Motion

- mass  $m$ , position  $x$ , velocity  $v$
- equations of motion:

$$\frac{d}{dt} x(t) = v(t)$$

$$\frac{d}{dt} v(t) = \frac{1}{m} F(x, v, t)$$

- Analytic solutions can be found for some classes of differential equations, but most can't be solved analytically
- Instead, we will numerically approximate a solution to our *initial value problem*

## Path Through a Field

- $f(\mathbf{X}, t)$  is a vector field defined everywhere
  - E.g. a velocity field which may change over time



- $\mathbf{X}(t)$  is a path through the field

## Higher Order ODEs

- Basic mechanics is a 2<sup>nd</sup> order ODE:

$$\frac{d^2}{dt^2} \mathbf{x} = \frac{1}{m} \mathbf{F}$$

- Express as 1<sup>st</sup> order ODE by defining  $\mathbf{v}(t)$ :

$$\frac{d}{dt} \mathbf{x}(t) = \mathbf{v}(t)$$

$$\frac{d}{dt} \mathbf{v}(t) = \frac{1}{m} \mathbf{F}(\mathbf{x}, \mathbf{v}, t)$$

$$\mathbf{X} = \begin{pmatrix} \mathbf{x} \\ \mathbf{v} \end{pmatrix} \quad f(\mathbf{X}, t) = \begin{pmatrix} \mathbf{v} \\ \frac{1}{m} \mathbf{F}(\mathbf{x}, \mathbf{v}, t) \end{pmatrix}$$

## For a Collection of 3D particles...

$$\mathbf{X} = \begin{pmatrix} p_x^{(1)} \\ p_y^{(1)} \\ p_z^{(1)} \\ v_x^{(1)} \\ v_y^{(1)} \\ v_z^{(1)} \\ p_x^{(2)} \\ p_y^{(2)} \\ p_z^{(2)} \\ v_x^{(2)} \\ v_y^{(2)} \\ v_z^{(2)} \\ \vdots \end{pmatrix} \quad f(\mathbf{X}, t) = \begin{pmatrix} v_x^{(1)} \\ v_y^{(1)} \\ v_z^{(1)} \\ \frac{1}{m_1} F_x^{(1)}(\mathbf{X}, t) \\ \frac{1}{m_1} F_y^{(1)}(\mathbf{X}, t) \\ \frac{1}{m_1} F_z^{(1)}(\mathbf{X}, t) \\ v_x^{(2)} \\ v_y^{(2)} \\ v_z^{(2)} \\ \frac{1}{m_2} F_x^{(2)}(\mathbf{X}, t) \\ \frac{1}{m_2} F_y^{(2)}(\mathbf{X}, t) \\ \frac{1}{m_2} F_z^{(2)}(\mathbf{X}, t) \\ \vdots \end{pmatrix}$$

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## Euler's Method

- Examine  $f(\mathbf{X}, t)$  at (or near) current state
- Take a step of size  $h$  to new value of  $\mathbf{X}$ :

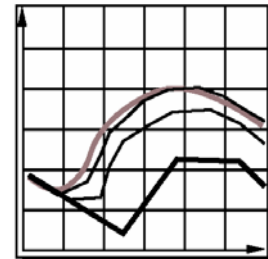
$$t_1 = t_0 + h$$

$$\mathbf{X}_1 = \mathbf{X}_0 + h f(\mathbf{X}_0, t_0)$$

- Piecewise-linear approximation to the curve

## Effect of Step Size

- Step size controls accuracy
- Smaller steps more closely follow curve
- For animation, we may want to take many small steps per frame



## Euler's Method: Inaccurate

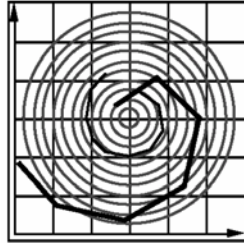
- Moves along tangent & can leave curve, e.g.:

$$f(\mathbf{X}, t) = \begin{pmatrix} -y \\ x \end{pmatrix}$$

- Exact solution is circle:

$$\mathbf{X}(t) = \begin{pmatrix} r \cos(t+k) \\ r \sin(t+k) \end{pmatrix}$$

- Euler's spirals outward no matter how small  $h$  is



## Euler's Method: Unstable

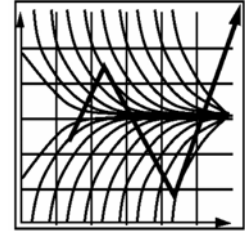
- Problem:  $f(x, t) = -kx$
- Solution:  $x(t) = x_0 e^{-kt}$

- Limited step size:

$$x_1 = x_0(1 - hk)$$

$$\begin{cases} h \leq 1/k & \text{ok} \\ h > 1/k & \text{oscillates } \pm \\ h > 2/k & \text{explodes} \end{cases}$$

- If  $k$  is big,  $h$  must be small



## Analysis using Taylor Series

- Expand exact solution  $\mathbf{X}(t)$

$$\mathbf{X}(t_0 + h) = \mathbf{X}(t_0) + h \left( \frac{d}{dt} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^2}{2!} \left( \frac{d^2}{dt^2} \mathbf{X}(t) \right) \Big|_{t_0} + \frac{h^3}{3!} (\dots) + \dots$$

- Euler's method:

$$\mathbf{X}(t_0 + h) = \mathbf{X}_0 + h f(\mathbf{X}_0, t_0) \dots + O(h^2) \text{ error}$$

$$h \rightarrow h/2 \Rightarrow \text{error} \rightarrow \text{error}/4 \text{ per step} \times \text{twice as many steps} \rightarrow \text{error}/2$$

- First-order method: Accuracy varies with  $h$ 
  - To get 100x better accuracy need 100x more steps

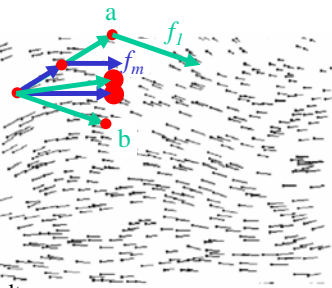
## Can we do better than Euler's Method?

- Problem:  $f$  has varied along the step
- Idea: look at  $f$  at the arrival of the step and compensate for variation



## 2nd-Order Methods

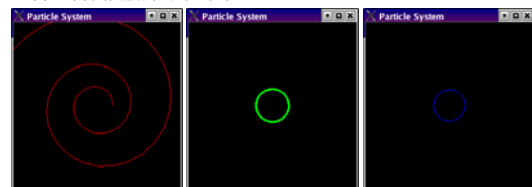
- Midpoint:**
  - 1/2 Euler step
  - evaluate  $f_m$
  - full step using  $f_m$
- Trapezoid:**
  - Euler step (a)
  - evaluate  $f_j$
  - full step using  $f_j$  (b)
  - average (a) and (b)
- Not exactly same result
- Same order of accuracy



## Comparison: Euler, Midpoint, Runge-Kutta

- initial position: (1,0,0)
- initial velocity: (0,5,0)
- force field: pulls particles to origin with magnitude proportional to distance from origin
- correct answer: circle

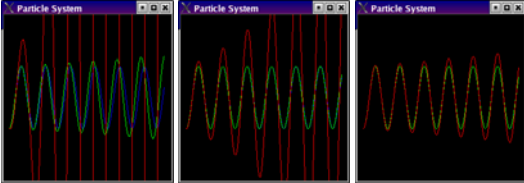
A 4<sup>th</sup> order method!



Euler will always diverge (even with small dt)

## Comparison: Euler, Midpoint, Runge-Kutta

- *initial position*: (0,-2,0) A 4<sup>th</sup> order method!
- *initial velocity*: (1,0,0)
- *force field*: pulls particles to line y=0 with magnitude proportional to distance from line
- *correct answer*: sine wave



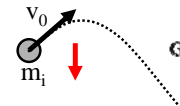
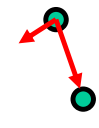
Decreasing the timestep (dt) improves the accuracy

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## Forces: Gravity

•For smoke, flame: make gravity point up!

- Simple gravity: depends only on particle mass
  - $\mathbf{v}_0$  
  - Gravity:  $\mathbf{f}^{(i)} = \begin{pmatrix} 0 \\ 0 \\ -m_i G \end{pmatrix}$
- N-body problem: depends on all other particles
  - Magnitude inversely proportional to square distance
  - $F_{ij} = G m_i m_j / r^2$  

## Forces: Spatial Fields

- Force on particle  $i$  depends only on position of  $i$ 
  - wind
  - attractors
  - repulsers
  - vortices
- Can depend on time
- Note: these add energy, may need damping too

## Forces: Damping

$$\mathbf{f}^{(i)} = -d\mathbf{v}^{(i)}$$

- Force on particle  $i$  depends only on velocity of  $i$
- Force opposes motion
- Removes energy, so system can settle
- Small amount of damping can stabilize solver
- Too much damping makes motion too glue-like

## Questions?

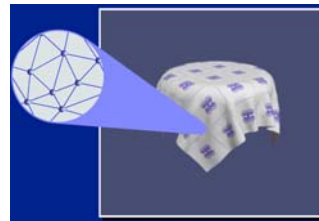


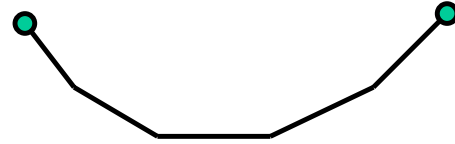
Image by Baraff, Witkin, Kass

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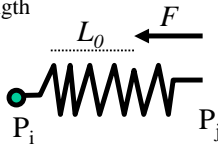
## How would you simulate a string?

- Each particle is linked to two particles
- Forces try to keep the distance between particles constant
- What force?



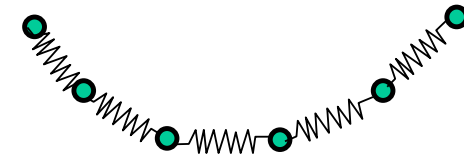
## Spring Forces

- Force in the direction of the spring and proportional to difference with rest length  $L_0$
- $$F(P_i, P_j) = K(L_0 - \|\vec{P}_i - \vec{P}_j\|) \frac{\vec{P}_i - \vec{P}_j}{\|\vec{P}_i - \vec{P}_j\|}$$
- $K$  is the stiffness of the spring
    - When  $K$  gets bigger, the spring really wants to keep its rest length



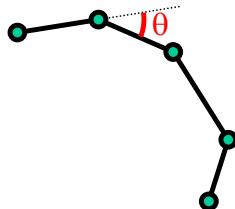
## How would you simulate a string?

- Springs link the particles
- Springs try to keep their rest lengths and preserve the length of the string
- Problems?
  - Stretch, actual length will be greater than rest length
  - Numerical oscillation



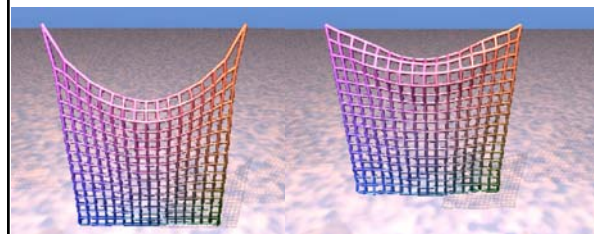
## How would you simulate hair?

- Similar to string...
- Deformation forces proportional to the angle between segments



## Reading for Today (2/6)

- “Deformation Constraints in a Mass-Spring Model to Describe Rigid Cloth Behavior”, Provat, 1995.

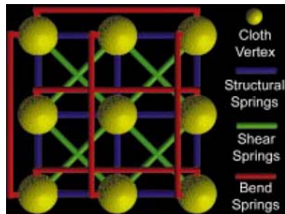


Simple mass-spring system

Improved solution

## Cloth Modeled with Mass-Spring

- Network of masses and springs
- Structural springs:
  - link  $(i, j)$  &  $(i+1, j)$   
and  $(i, j)$  &  $(i, j+1)$
- Shear springs
  - link  $(i, j)$  &  $(i+1, j+1)$   
and  $(i+1, j)$  &  $(i, j+1)$
- Flexion (Bend) springs
  - link  $(i, j)$  &  $(i+2, j)$   
and  $(i, j)$  &  $(i, j+2)$



From Lander

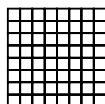
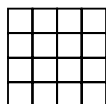
<http://www.darwin3d.com/gamesdev/articles/col0599.pdf>

## The Stiffness Issue

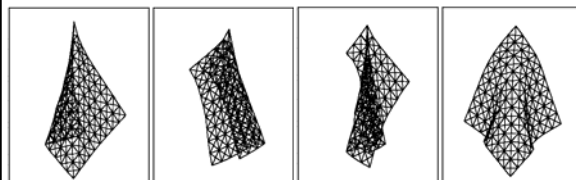
- What relative stiffness do we want for the different springs in the network?
- Cloth is barely elastic, shouldn't stretch so much!
- Inverse relationship between stiffness &  $\Delta t$
- We really want a constraints (not springs)
- Many numerical solutions
  - reduce  $\Delta t$
  - use constraints
  - implicit integration
  - ...

## The Discretization Problem

- What happens if we discretize our cloth more finely, or with a different mesh structure?
- Do we get the same behavior?
- Usually not! It takes a lot of effort to design a scheme that does not depend on the discretization.



## Questions?



*Interactive Animation of  
Structured Deformable Objects*  
Desbrun, Schröder, & Barr 1999

## Reading for Tuesday 2/13:

- “Realistic Animation of Liquids”,  
Foster & Metaxas, 1996

