27. Impact Mechanics of Manipulation

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Outline.

- Impulse-momentum equations
- Impact of a particle
 - Plastic and elastic impact
 - Newtonian and Poisson restitution
 - Impulse space
- Impact of a planar rigid body
 - Solution of the sliding rod problem

Frictional impact

When rigid bodies collide:

- Discontinuity in velocity,
- Infinite forces,
- Zero time,
- Even for simple cases there is no agreement on laws,
- For multiple contact indeterminacies abound.

We will adopt Routh's analysis of planar impact, which uses Poisson restitution.

Impulse momentum equations

Recall:

$$\mathbf{F} = \frac{d\mathbf{P}}{d\mathbf{t}}$$
$$\mathbf{N} = \frac{d\mathbf{L}}{d\mathbf{t}}$$

Integrating:

$$\Delta \mathbf{P} = \int_{t_0}^{t_1} \mathbf{F} \, dt$$
$$\Delta \mathbf{L} = \int_{t_0}^{t_1} \mathbf{N} \, dt$$

where $\Delta \mathbf{P} = \mathbf{P}_1 - \mathbf{P}_0$, $\Delta \mathbf{L} = \mathbf{L}_1 - \mathbf{L}_0$, $\int \mathbf{F} dt$ is the *impulse*, and $\int \mathbf{N} dt$ is the *impulsive moment*.

Determining impulse. Particle in plane.

We need impulse and impulsive moment as a function of state and physical parameters.

Start with a frictionless particle on plane. Velocity changes discontinuously from v_0 to v_1 :

$$\Delta \mathbf{v} = \frac{1}{m} \Delta \mathbf{P} = \frac{1}{m} \mathbf{I}$$

where

$$\Delta \mathbf{v} = \mathbf{v}_1 - \mathbf{v}_0$$

The tangential impulse I_t is zero, so

$$v_{1t} = v_{0t}$$

 $v_{1n} = v_{0n} + \frac{1}{m}I_n$

Plastic and elastic

How to get the normal impulse I_n ? Two special cases: *plastic impact*

$$I_n = -mv_{0n}$$
$$\rightarrow v_{1n} = 0$$

and *elastic impact*

 $I_n = -2mv_{0n}$ $\rightarrow v_{1n} = -v_{0n}$

Newtonian restitution

Newton hypothesized a continuum from plastic to elastic, defined by a *coefficient of restitution*

$$e = -\frac{v_{1n}}{v_{0n}} \tag{1}$$

so that plastic impact corresponds to e = 0, and elastic impact corresponds to e = 1.

Poisson restitution

A competing definition of the coefficient of restitution was given by Poisson. Divide the collision into two stages.

 $v_n < 0$ compression $v_n = 0$ $v_n > 0$ restitution

Let I_c be impulse accumulated during compression, and let I_r be impulse accumulated during restitution. Poisson's hypothesis is that the ratio of these two parts is governed by material properties:

$$e = \frac{I_r}{I_c} \tag{2}$$

For simple cases no advantage over Newton, but for frictional rigid bodies Poisson restitution is superior.

Applying Poisson restitution for frictionless parti

Given m, \mathbf{v}_0 , and e, first find I_c :

$$I_c = m v_{0n}$$

then restitution impulse I_r

$$I_r = emv_{0n}$$

then total impulse I

$$I = I_c + I_r = (1+e)mv_{0n}$$

and resulting normal velocity:

$$v_{1n} = -ev_{0n}$$

The tangential velocity is unchanged.

Frictional particle

$$\hat{\mathbf{n}} \bigvee \mathbf{v}_0 = \begin{pmatrix} -\epsilon \\ -1 \end{pmatrix} \\ \mathbf{v}_1 = \begin{pmatrix} 3 - \epsilon \\ 0 \end{pmatrix} \\ \hat{\mathbf{t}} \end{pmatrix}$$

Consider frictional particle with initial leftward velocity. There should be some frictional impulse acting toward the right. Suppose we apply Coulomb's law, by looking at initial contact mode, and multiplying I_n by μ . The result is absurd, and violates conservation of energy!

A better model



The left-drifting particle



Rigid body. Kinematics.

First, we observe some kinematic relations

$$\mathbf{c} = \mathbf{x} + \mathbf{r}$$
$$\dot{\mathbf{c}} = \dot{\mathbf{x}} + \dot{\mathbf{r}}$$
$$= \dot{\mathbf{x}} + \boldsymbol{\omega} \times \mathbf{r}$$
$$\Delta \dot{\mathbf{c}} = \Delta \dot{\mathbf{x}} + \Delta \boldsymbol{\omega} \times \mathbf{r}$$



Impulse momentum laws

Now we can write the impulse-momentum laws:

$$m\Delta \dot{\mathbf{x}} = \mathbf{I}$$
 (3)

$$\rho^2 m \Delta \omega = \mathbf{r} \times \mathbf{I} \tag{4}$$

Substituting into the expression for $\Delta \dot{\mathbf{c}}$ yields:

$$\Delta \dot{\mathbf{c}} = \frac{1}{m} \mathbf{I} + \frac{1}{\rho^2 m} (\mathbf{r} \times \mathbf{I}) \times \mathbf{r}$$
(5)

$$=\frac{1}{m}\mathbf{I} - \frac{1}{\rho^2 m}\mathbf{r} \times (\mathbf{r} \times \mathbf{I})$$
(6)

$$=\frac{1}{\rho^2 m} \left(\rho^2 I_3 - R^2\right) \mathbf{I} \tag{7}$$

where I_3 is the three by three identity matrix, and R is the cross-product matrix.

Lecture 27.

Compression line and restitution line

Substituting above and expanding, we obtain

$$\begin{pmatrix} \Delta \dot{c}_t \\ \Delta \dot{c}_n \end{pmatrix} = \frac{1}{\rho^2 m} \begin{pmatrix} (\rho^2 + r_n^2) & -r_t r_n \\ -r_t r_n & (\rho^2 + r_t^2) \end{pmatrix} \begin{pmatrix} I_t \\ I_n \end{pmatrix}$$
(8)

The sticking line is defined by the equation $\dot{c}_t = 0$,

$$\dot{c}_{0t} + \frac{\rho^2 + r_n^2}{\rho^2 m} I_t - \frac{r_t r_n}{\rho^2 m} I_n = 0$$
(9)

and the compression line is defined the equation $\dot{c}_n = 0$,

$$\dot{c}_{0n} - \frac{r_t r_n}{\rho^2 m} I_t + \frac{\rho^2 + r_t^2}{\rho^2 m} I_n = 0$$
(10)

Rigid body example



Analysis of rigid body example



Sliding rod problem



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Sliding rod solution



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